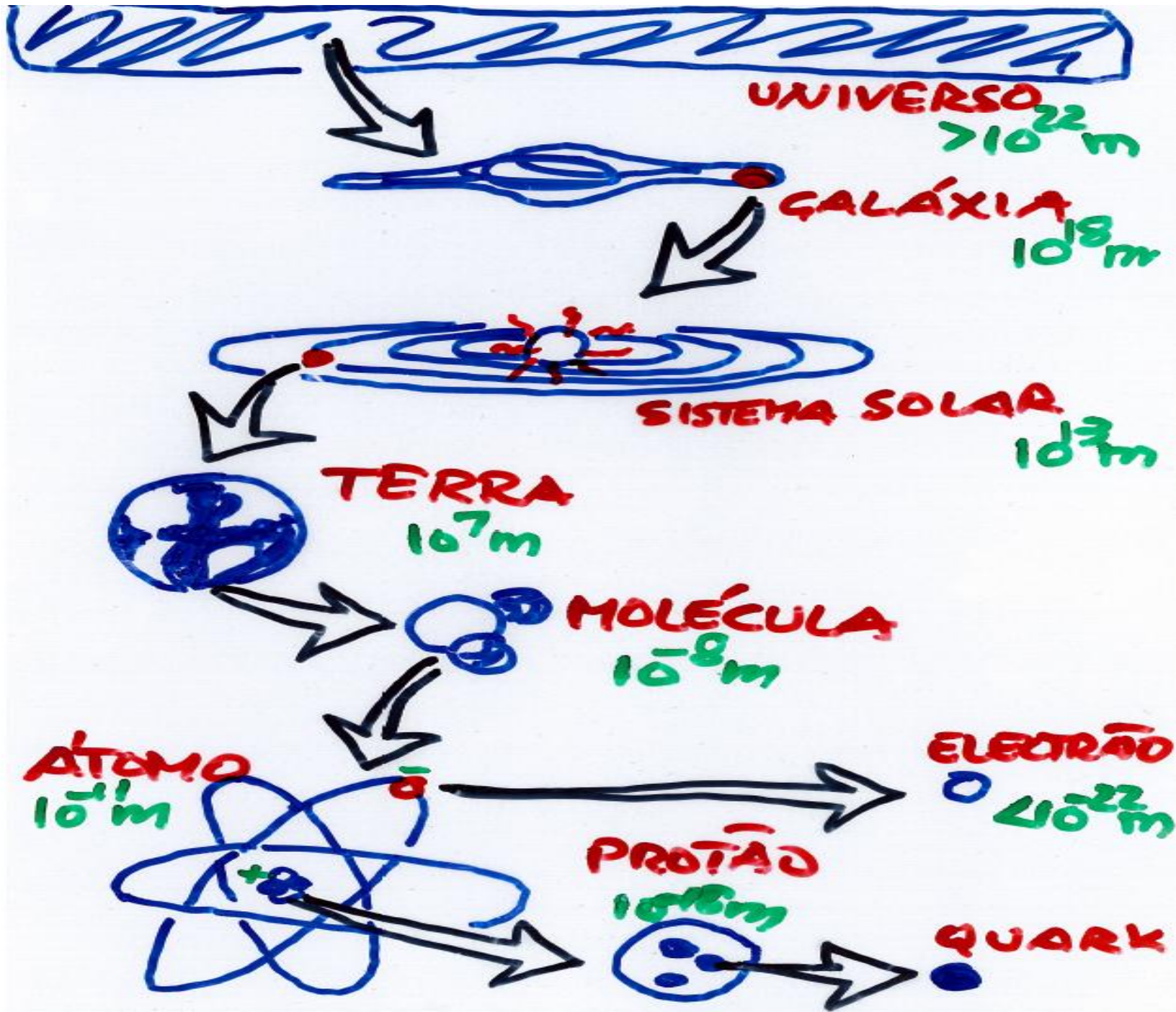


INTERAÇÕES  
FUNDAMENTAIS  
NO  
UNIVERSO

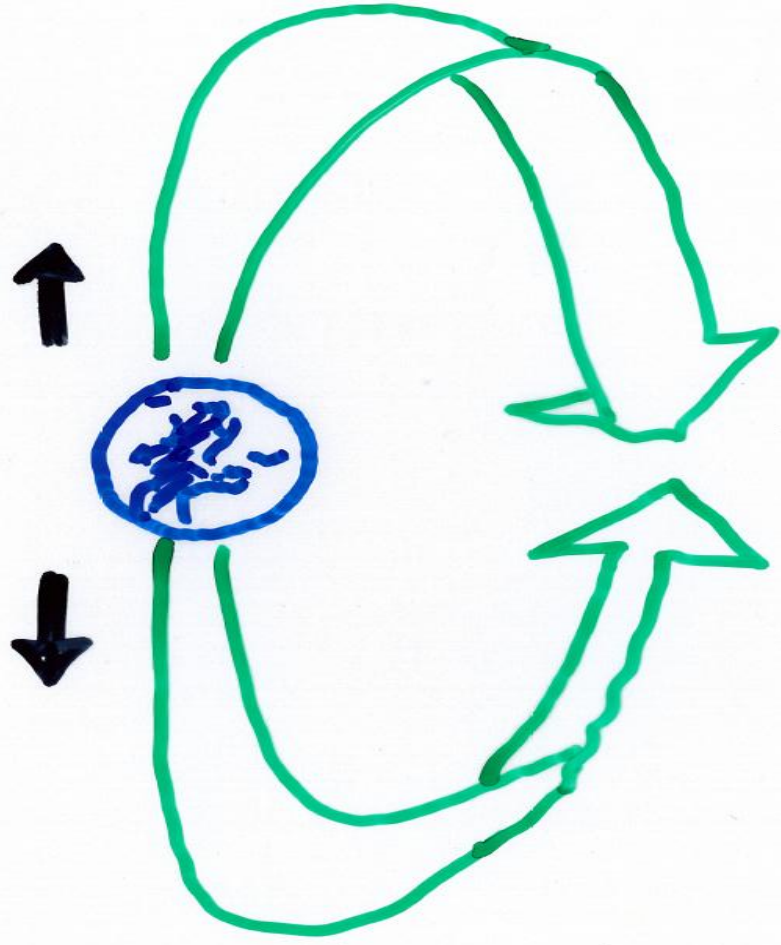
J. DAS DEUS

Ea9 2012



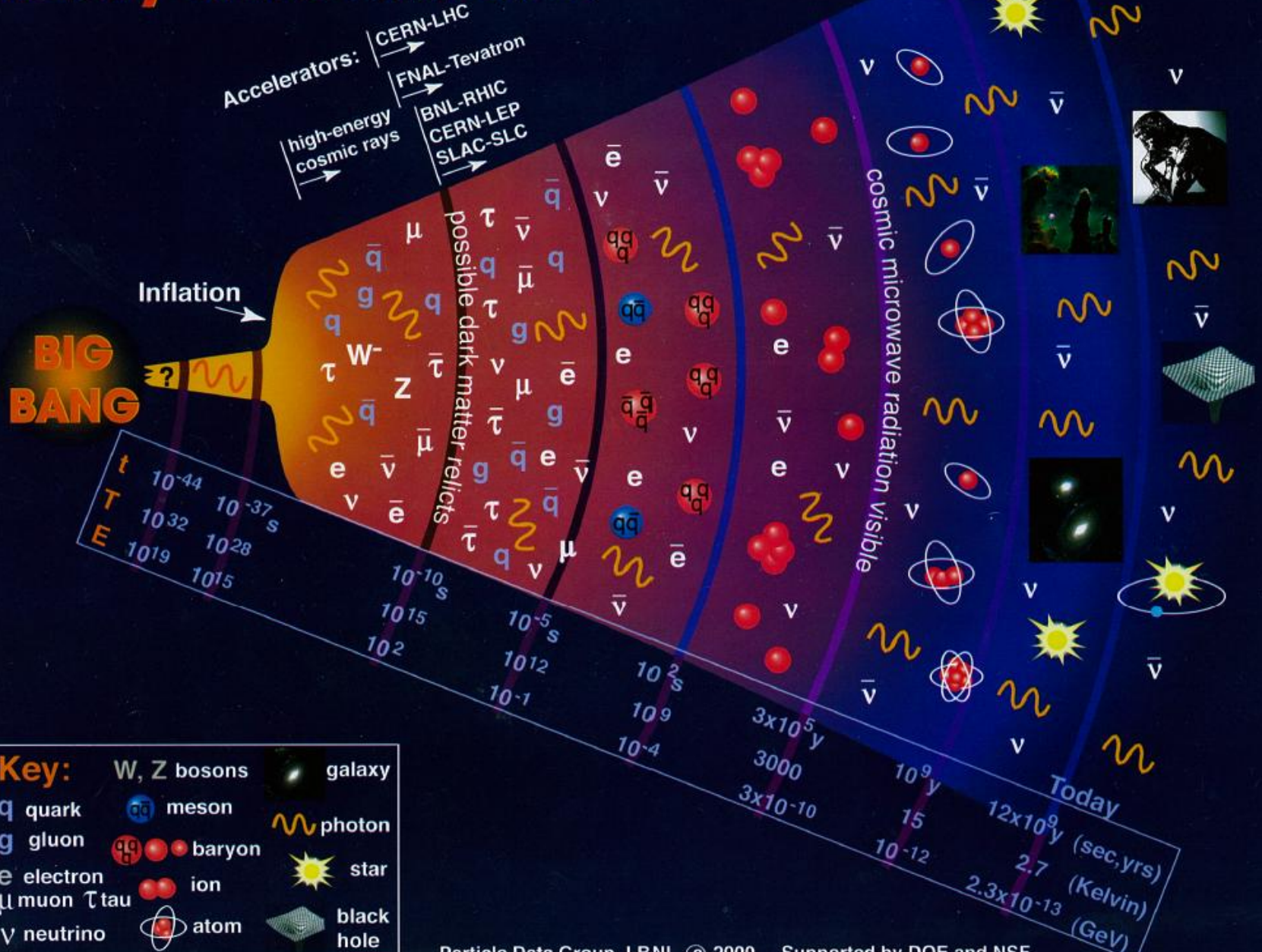
ACELERADORES

TELESCÓPIOS





# History of the Universe



# OBJECTOS + INTERACÇÕES

$$\vec{P} = \text{const.} \Rightarrow \frac{d\vec{P}}{dt} = \vec{0}$$

$$\frac{d\vec{P}}{dt} = \vec{F}$$

$$\frac{d\vec{L}}{dt} = \vec{r} \times \vec{F} \equiv \vec{N}$$

$$\frac{dE}{dt} = \vec{F} \cdot \vec{v}$$



TRÊS INVARIÂNCIAS:

TRANSLAÇÃO NO ESPAÇO

ROTAÇÃO NO ESPAÇO

TRANSLAÇÃO NO TEMPO

TRÊS LEIS DE CONSERVAÇÃO

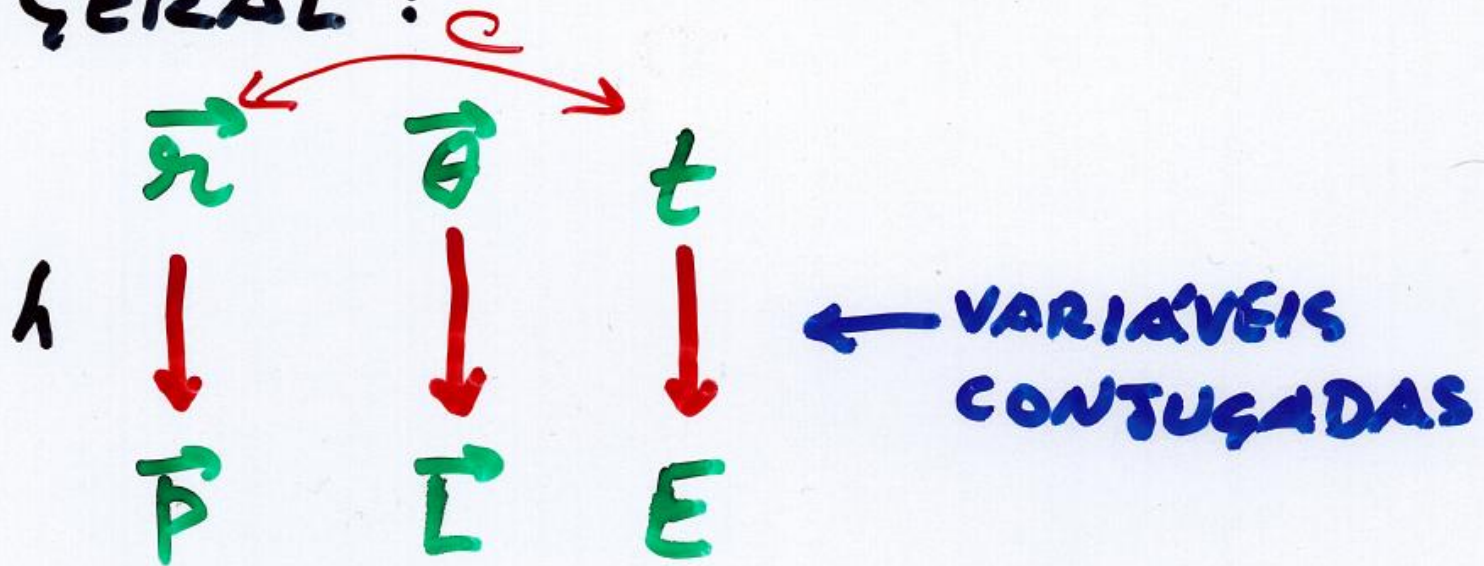
$$d\vec{P}/dt = 0$$

$$d\vec{L}/dt = 0$$

$$dE/dt = 0$$

Noether

EM GERAL :



$$(x, y, z, t) \longrightarrow (p_x, p_y, p_z, E)$$

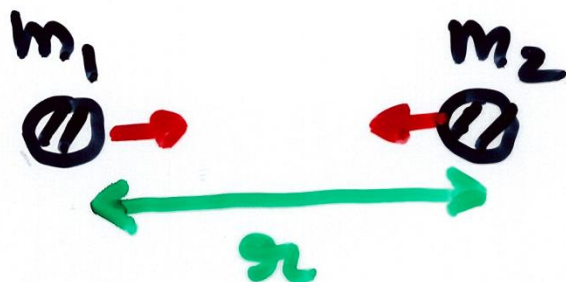
NOTA :

→ LEIS DA FÍSICA

→ MECÂNICA QUÂNTICA

# CARGAS & CAMPOS

Newton :



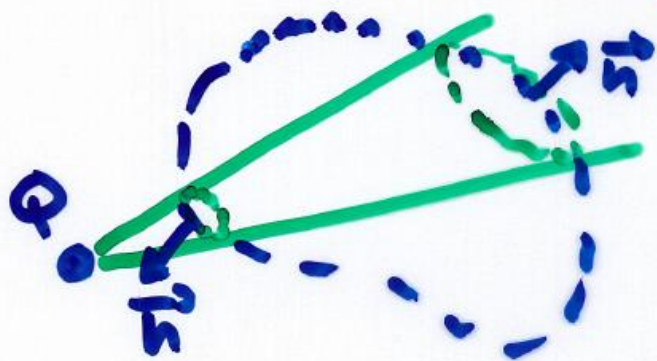
$$|\vec{F}| = G \frac{m_1 m_2}{r^2}$$

{	$m_1$	→	$Q$	CARÇA	
	$\text{⊗}$	→		CAMPO	$\vec{E}_1 = -G \frac{m_1}{r^2} \rightarrow -\frac{q Q_1}{r^2}$
	$G$	→	$\alpha$		$F = Q_2 E_1$



# TEOREMA DE GAUSS

$$\Phi = \int_{\text{[Fechada]}} \vec{E} \cdot d\vec{S} = 4\pi\alpha Q|_{\text{interior}}$$



$$\Phi = 0$$



$$\Phi > 0$$

Newton :  
(Coulomb



$$\bar{\Phi} = \int \frac{\alpha Q}{r^2} dS = \frac{\alpha Q}{r^2} 4\pi r^2 = 4\pi \alpha Q$$

Em geral:

$$S \rightarrow r^{D-1}$$

$$r^2 \rightarrow r^{\tilde{\alpha}}$$

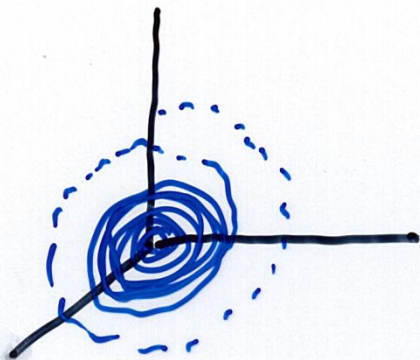
T. Gauss

$$\frac{r^{D-1}}{r^{\tilde{\alpha}}} \Rightarrow \boxed{\tilde{\alpha} = D-1}$$

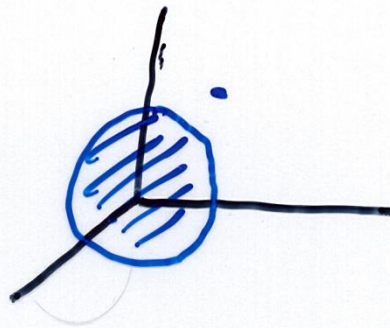
D=1:



# ○ MISTÉRIO DA MATÉRIA ESCURA :

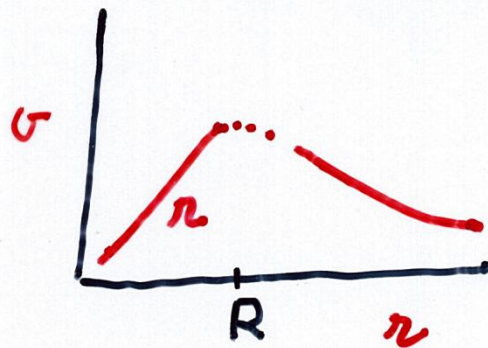
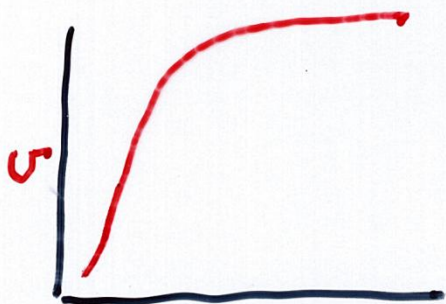


~



$$G \frac{M m}{r^2} = \frac{m v^2}{r} \quad \left\{ \begin{array}{l} M \sim r^3, \quad r \leq R \\ M \sim R^3, \quad r \geq R \end{array} \right.$$

$$\left\{ \begin{array}{l} v \sim r \\ v \sim 1/r^{1/2} \end{array} \right.$$



EXP.

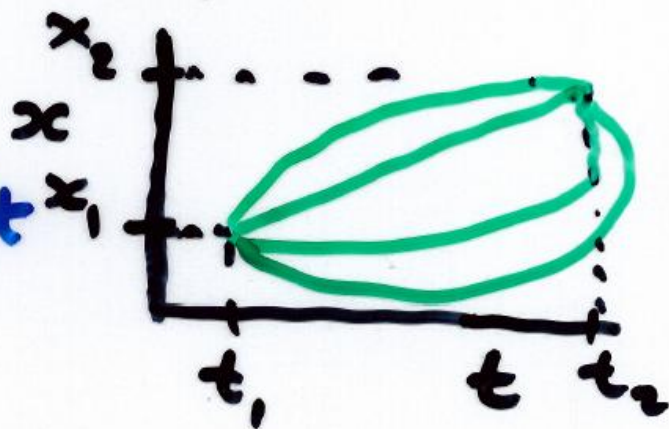
Zwicky



# COMO DESCREVER TEORIAS DE CAMPO

LAGRANGEANO  $\rightarrow$  AÇÃO

CLÁSSICO:  $L(x, \dot{x}) = \frac{1}{2} m \dot{x}^2 - U(x)$

$$S[x(t)] = \int_{t_1}^{t_2} L(x(t), \dot{x}(t)) dt$$


Hamilton

EQ. Euler-  
Lagrange

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{x}} - \frac{\partial L}{\partial x} = 0$$

NOTA: SISTEMA ISOLADO

$$\frac{\partial L}{\partial t} = 0; \quad \frac{\partial L}{\partial x} = 0; \quad \frac{\partial L}{\partial \theta} = 0$$

---

PARTÍCULA  $\longrightarrow$  CAMPO

$x$  (•)  $\longrightarrow \phi(x)$  

$$\text{Ex: } \mathcal{L} = -\frac{1}{2}(\partial_\mu \phi)^2 - \frac{1}{2}m^2\phi^2 - \lambda\phi^3 - g\phi^4$$

$$S[\phi(x)] = \int d^4x \mathcal{L}(\phi(x), \partial_\mu \phi(x))$$

$$\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \right) - \frac{\partial \mathcal{L}}{\partial \phi} = 0$$



# EX: QED

ELECTRÕES

$$\mathcal{L}_\psi = -\bar{\psi} \partial_\mu \psi - m \bar{\psi} \psi$$

$$[\bar{\psi} \equiv \psi^*]$$

INVARIÂNCIAS (PADRÃO)

GLOBAL:  $\psi \rightarrow \psi' = e^{i q \epsilon} \psi$  ✓

LOCAL:  $\psi \rightarrow \psi' = e^{i q \epsilon(x)} \psi$  ?

$$\partial_\mu \psi(x) \rightarrow (\partial_\mu \psi(x))' = e^{i q \epsilon(x)} [\partial_\mu \psi + i q \partial_\mu \epsilon(x) \psi]$$

NÃO É INVARIANTE

DERIVADA COVARIANTE  $D_\mu$ :

$$D_\mu \psi(x) \rightarrow (D_\mu \psi(x))' = e^{i q \epsilon(x)} (D_\mu \psi)$$

$$\Rightarrow D_\mu \rightarrow \partial_\mu - i q A_\mu(x)$$

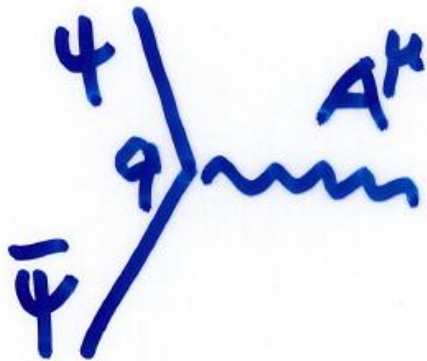


$$A_\mu \rightarrow A'_\mu = A_\mu + \partial_\mu \epsilon(x) \quad \text{CAMPO POTENCIAL (GAUGE)}$$

$$\mathcal{L}_\psi = -\bar{\psi} \not{\partial} \psi - m \bar{\psi} \psi + i q A_\mu \bar{\psi} \gamma^\mu \psi$$

$$\mathcal{L}_{A+\psi} = -\frac{1}{4} (\partial_\mu A_\nu - \partial_\nu A_\mu)^2 - \bar{\psi} \not{\partial} \psi - m \bar{\psi} \psi + i q A_\mu \bar{\psi} \gamma^\mu \psi$$

[Conservação da corrente]



# INVARIÂNCIA PADRÃO GLOBAL

$$\uparrow \equiv \phi$$



# INVARIÂNCIA PADRÃO LOCAL

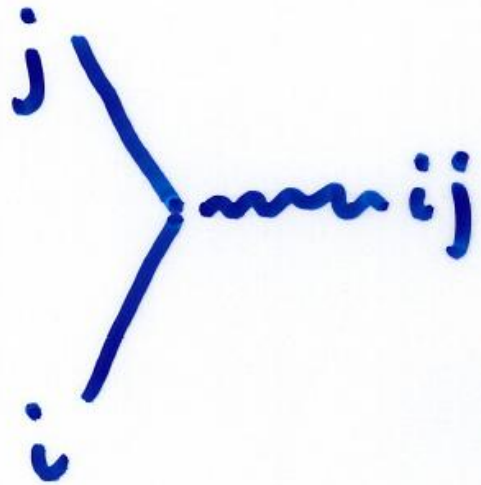
$$\uparrow \equiv A_\mu$$



# GENERALIZAÇÕES

1)  $\Psi_i(x) \rightarrow \Psi'_i(x) = U_{ij} \Psi_j(x)$

$$U^\dagger = U^{-1}$$



$$SU(N): N \otimes N = (N^2 - 1) \oplus 1$$

[N]: FUNDAMENTAL  $S = \frac{1}{2}$

$[N^2 - 1]$ : GERADORES  $S = 1$

[Yang-Mills]



# OBSERVAÇÕES:

1) OS FOTÕES NÃO INTERAÇEM ENTRE SI !



$$1[2,3] - 1[3,2]$$

$U(1)$	GR. ABELIANO	$[a,b] - [b,a] = 0$
$SU(N)$	GR. NÃO ABELIANO	$[a,b] - [b,a] \neq 0$



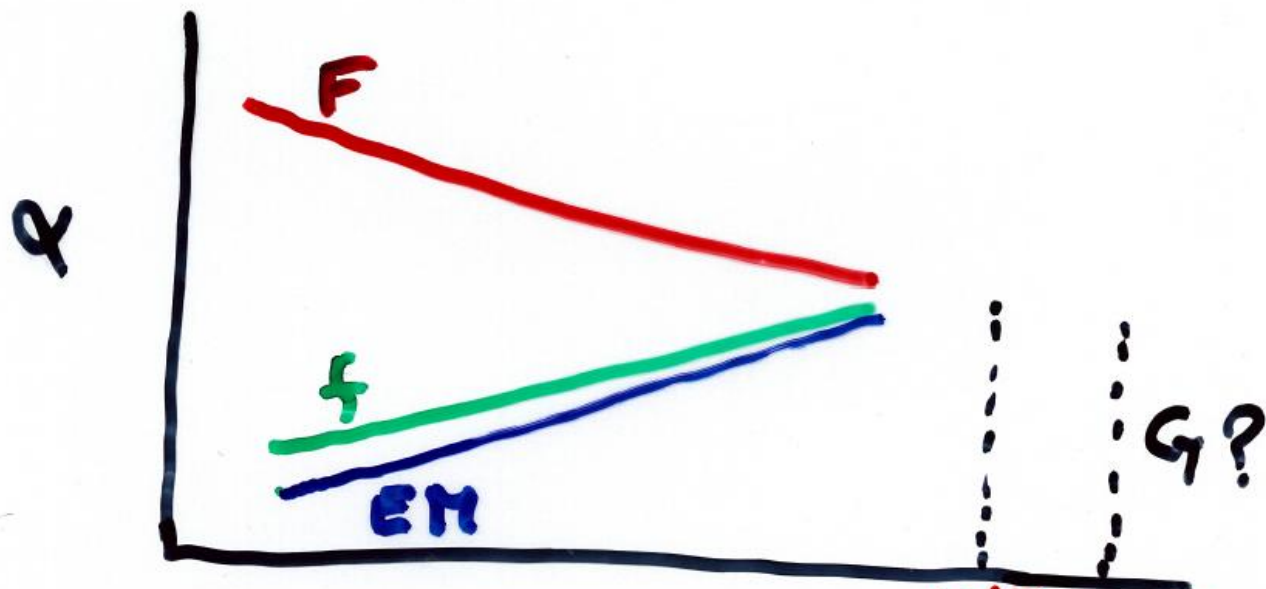
# 27 GUTs

EM  
If  
IF

$$\alpha \equiv \frac{1}{4\pi\epsilon_0} \frac{e^2}{\hbar c}$$

$g_f$   
 $g_F$

sem dimensões



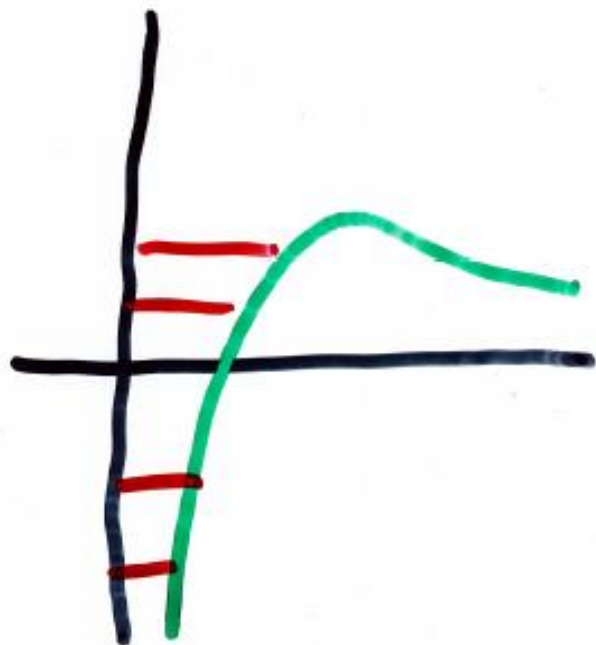
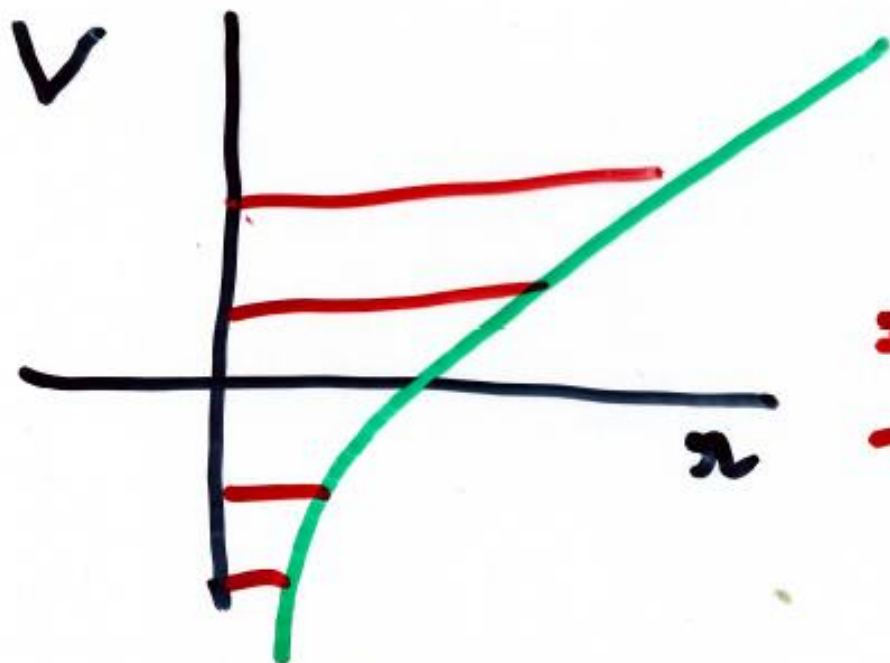
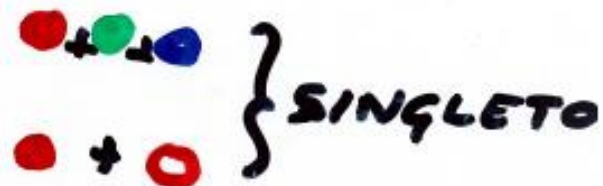
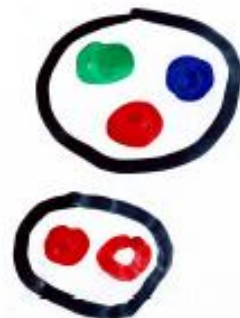
[ALCANÇE?]

ENERGIAS  $10^{16}$   $10^{19}$  GeV

### 3) CONFINAMENTO

NUCLEÕES

MESÕES



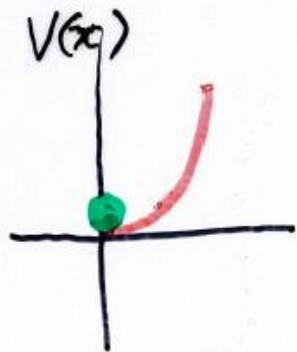
NOTA: PRIMEIRAS 4-SEGUNDAS...



# QUEBRA ESPONTÂNEA DE SIMETRIA (LANDAU)

EX: ATOMO DE HIDROGÊNIO  
SUPER CONDUCTOR  
FERRO MAGNETO

$$V(x) = -\frac{1}{2}\mu^2 x^2 + \frac{1}{4}\lambda x^4 = \left(-\frac{1}{2}\mu^2 + \frac{1}{4}\lambda x^2\right)x^2$$



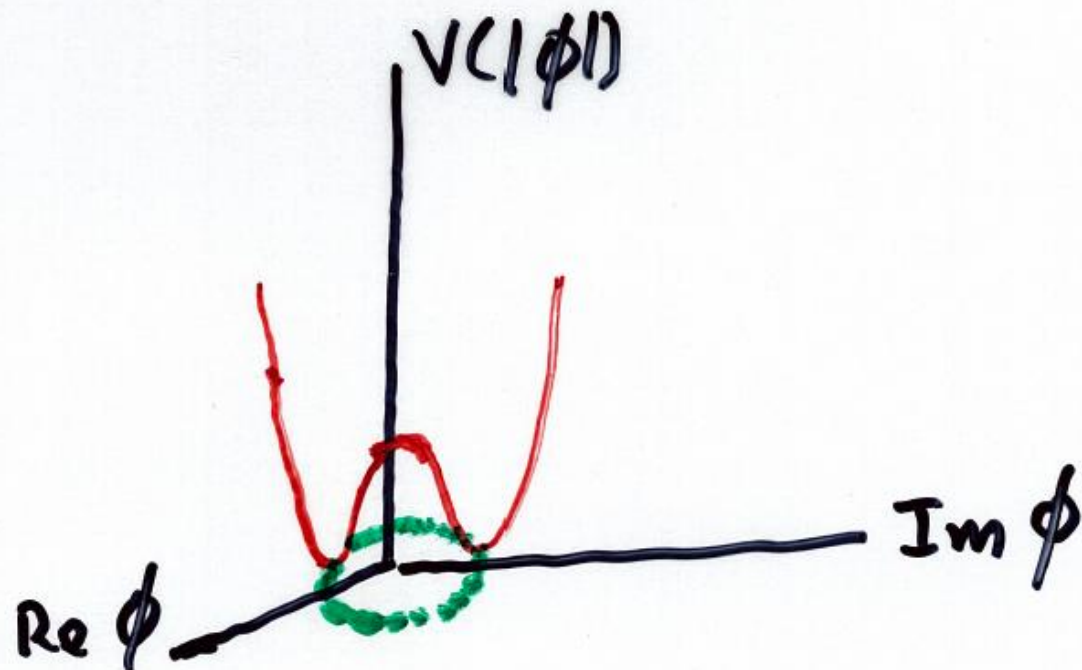
$$\mu^2 < 0$$



$$\mu^2 > 0$$

→ QUEBRA ESPONTÂNEA DE SIMETRIA:  
QUANDO A SIMETRIA DA TEORIA  
NÃO APARECE REFLECTIDA NA  
SOLUÇÃO

EX: INVARIÂNCIA DE ROTAÇÃO



CAMPO ESCALAR COMPLEXO  $[\text{Re}\phi, \text{Im}\phi]$

LAGRANGIANO:

$$\mathcal{L} = -\partial_\mu |\phi|^2 - V(|\phi|)$$

$$\phi(x) \longrightarrow \phi'(x) = e^{iq\epsilon} \phi(x) \quad (\text{INV. PADRÃO GLOBAL})$$

SOLUÇÕES: EM TORNO DO MÍNIMO DE POTENCIAL

$$\text{CASO "NORMAL": } V(|\phi|) = V(|\phi|)_{\text{MIN.}} + \mu^2 |\phi|^2 + \dots$$

$$V(|\phi|)_{\text{MIN.}} = 0$$

$$\mu^2 \equiv \frac{\partial V}{\partial |\phi|^2} \equiv \text{massa}$$



# CASO "QUEBRA DE SIMETRIA"



$$\phi(x) = \frac{1}{\sqrt{2}} \rho(x) e^{i\theta(x)}$$

$$\partial_\mu \phi = \frac{1}{\sqrt{2}} e^{i\theta(x)} (\partial_\mu \rho + i \rho \partial_\mu \theta)$$

$$\mathcal{L} = -\frac{1}{2} (\partial_\mu \rho)^2 - \frac{1}{2} \rho^2 (\partial_\mu \theta)^2 - V(\rho/\sqrt{2})$$

EXP. EM TORNO MÍNIMO:

$$\rho : \mu^2 = \left. \frac{\partial^2 V(\rho/\sqrt{2})}{\partial \rho^2} \right|_{\rho_{\text{MIN}}}$$

$\theta$  : massa zero  
(GOLDSTONE)

# INVARIÂNCIA GAUSSIANA LOCAL <sup>4</sup>

(Brout-Englert-Higgs)

$$\phi(x) \rightarrow \phi'(x) = e^{iqE(x)} \phi(x)$$

$$A_\mu(x) \rightarrow A'_\mu(x) = A_\mu(x) + \partial_\mu E(x)$$

⇒

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^2(A) - |D_\mu \phi|^2 - V(|\phi|)$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$D_\mu \phi = \partial_\mu \phi - iqA_\mu \phi$$

POTENCIAL:  $\Theta(x) \rightarrow \Theta'(x) = \Theta(x) + qE(x)$

DÉRIVADA CO-VARIANTE:

$$D_\mu \phi = \frac{1}{\sqrt{2}} e^{i\theta} (\partial_\mu \rho - iq\rho (A_\mu - g' \partial_\mu \theta))$$

$B_\mu$  (insensível à transformação de gauge)

$$D_\mu \phi = \frac{1}{\sqrt{2}} e^{i\theta} (\partial_\mu \rho - iq\rho B_\mu)$$

$$\Rightarrow \mathcal{L} = -\frac{1}{4} F_{\mu\nu}^2(B) - \frac{1}{2} (\partial_\mu \rho)^2 - \frac{1}{2} q^2 \rho^2 B_\mu^2 - V(\rho/\sqrt{2})$$

$$M_B = |q\rho|_{\text{MIN}} = |q\upsilon|$$

MASSA DE HIGGS

$$V(\rho) = -\frac{1}{2} \mu^2 \rho^2 + \frac{1}{4} \lambda \rho^4$$

$$\text{MIN: } -\mu^2 + \lambda \rho_{\text{MIN}}^2 = 0 \longrightarrow \upsilon = \sqrt{\mu^2/\lambda}$$

$$M_{\text{HIGGS}}^2 = 2\lambda \upsilon^2$$