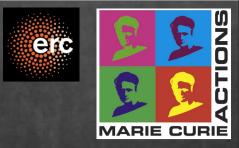


6ª Escola de Astrofísica e Gravitação do IST Lisboa – Septembro 4-8, 2012



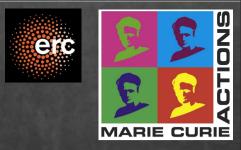
Ondas Gravitacionais

 $g_{\mu\nu} \sim \eta_{\mu\nu} + h_{\mu\nu}$ $\Box \bar{h}_{ab} = -16\pi G T_{ab}$

Paolo Pani CENTRA – Instituto Superior Técnico <u>http://blackholes.ist.utl.pt</u>



6th School in Astrophysics and Gravitation <u>IST- Lisboa – September 4-8, 2012</u>



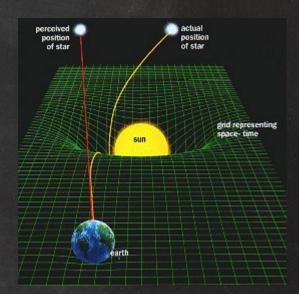
Gravitational waves

 $g_{\mu\nu} \sim \eta_{\mu\nu} + h_{\mu\nu}$ $\Box \bar{h}_{ab} = -16\pi G T_{ab}$

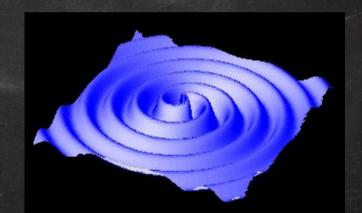
Paolo Pani CENTRA – Instituto Superior Técnico <u>http://blackholes.ist.utl.pt</u>

When, how, why

- Predicted by Einstein in 1916, right after GR
- Ripples in spacetime
- Cannot exist in Newtonian gravity
- Deformations from spherical symmetry
- GWs are weak! (pro and cons)

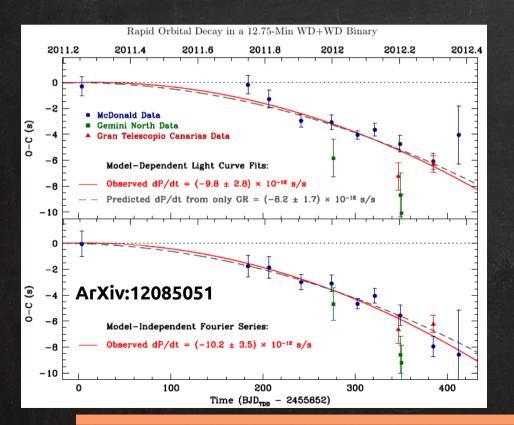


Light deflection in GR



Do GWs exist?

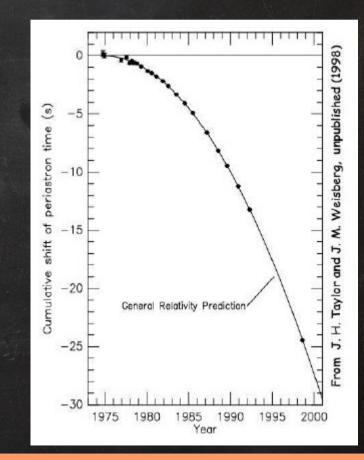
- No direct evidence yet
- Indirect evidences
- Hulse-Taylor pulsar PSR B1913+16
- WD+WD binary J0651+2844D





Nobel prize 1993

Weber's bar

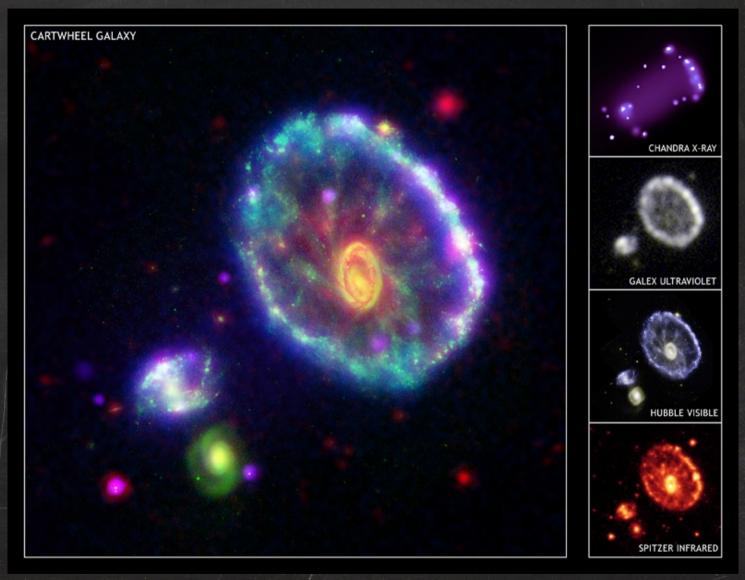


GW (astro)physics

- GW astronomy
- Tests of GR
- Big bang echoes
- Gamma-ray bursts
- Accretion disks
- Galaxy populations

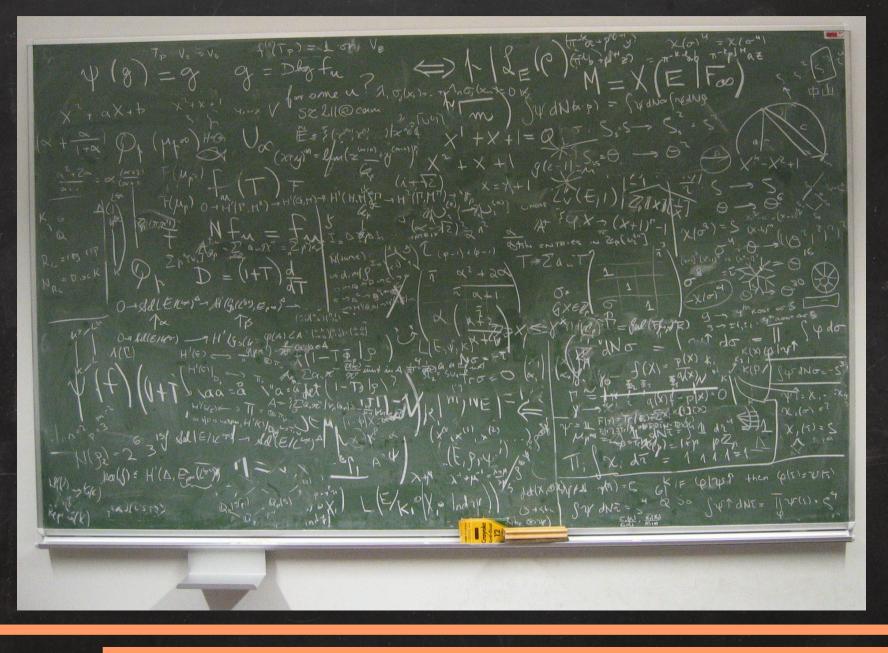
Credit: MPI for Gravitational Physics/W.Benger-Z

All eyes (...and ears)



Credit: Chandra, GALEX, Hubble, Spitzer -NASA/JPL/Caltech/P.Appleton et al.

A lot of theory



Linearized GR

$$R_{ab} - \frac{1}{2}g_{ab}R + \Delta g_{ab} = 8\pi G T_{ab}$$

 $g_{ab} = \eta_{ab} + \epsilon h_{ab} + \mathcal{O}(\epsilon^2)$

$$\Gamma^a_{bc} = \epsilon \frac{\eta^{al}}{2} (h_{lc,b} + h_{lb,c} - h_{bc,a})$$

$$R_{abcd} = \frac{\epsilon}{2} (h_{ad,bc} + h_{bc,ad} - h_{ac,db} - h_{bd,ac})$$

$$R_{ab} = \eta^{cd} R_{cadb} = \frac{\epsilon}{2} (h_{a,bc}^c + h_{b,ac}^c - \Box h_{ab} - h_{,ab})$$
$$R = \eta^{ab} R_{ab} = \epsilon (h_{,cd}^{cd} - \Box h)$$

$$G_{ab} = \frac{\epsilon}{2} (h_{a,bc}^{c} + h_{b,ac}^{c} - \Box h_{ab} - h_{,ab} - \eta_{ab} h_{,cd}^{cd} + \eta_{ab} \Box h)$$

Gauge freedom

$$x^a \to {x'}^a = x^a + \epsilon \xi^a$$

$$h'_{ab} = h_{ab} - 2\epsilon \xi_{(b,a)}$$

Curvature tensors are gauge invariant

$$\bar{h}_{ab} \equiv h_{ab} - \frac{1}{2}\eta_{ab}h$$

$$G_{ab} = \frac{\epsilon}{2} (\bar{h}^c_{a,bc} + \bar{h}^c_{b,ac} - \Box \bar{h}_{ab} - \eta_{ab} \bar{h}^{cd}_{,cd})$$

Transverse trace-reserved gauge: $\bar{h}^a_{b.a} = 0$ $\Box \bar{h}_{ab} = -16\pi G T_{ab}$

In vacuum:

 $\Box h_{ab} = 0 \Rightarrow \Box R_{abcd} = 0$

Degrees of freedom

Irreducible decomposition:

$$h_{ab} = \begin{pmatrix} -2\Phi & w_i \\ w_i & h_{ij} \end{pmatrix} \qquad h_{ij} = 2s_{ij} - 2\Psi\delta_{ij}$$

$$\Box s_{ij} = 0 \Rightarrow \Box h_{ij}^{TT} = 0 \qquad h_{ij}^{TT} = C_{ij} e^{ikx}$$
 Plane wave

Geodesic deviation equation:

$$\frac{D^2}{d\tau^2}S^{\mu} = R^{\mu}_{\nu\rho\sigma}u^{\nu}u^{\rho}S^{\sigma} \qquad \qquad \frac{D}{d\lambda} = \frac{dx^{\mu}}{d\lambda}\nabla_{\mu}$$

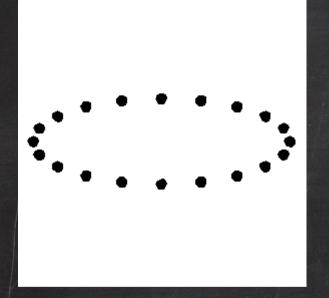
$$R_{\mu 00\sigma} = \frac{1}{2} \partial_0 \partial_0 h_{\mu\sigma}^{TT}$$

$$S^{1,2}(t) = (1 \pm \frac{1}{2}h_+e^{ikx})S^{1,2}(0)$$

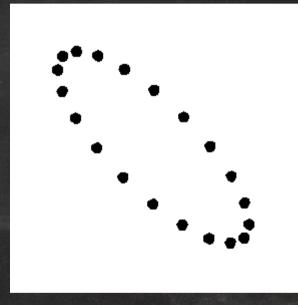
Plus (+) polarization

$$S^{1,2}(t) = S^{1,2}(0) + \frac{1}{2}h_x e^{ikx} S^{2,1}(0)$$

Degrees of freedom







Cross (x) polarization

~ lowest-energy modes of a string

Production of GWs (1)

 $\Box \bar{h}_{ab} = -16\pi G T_{ab}$

Green function approach

 $\Box G(x-y) = \delta^{(4)}(x-y) \quad \bar{h}_{ab}(x) = -16\pi G \int d^4 y G(x-y) T_{ab}(y)$

Retarded Green function

$$G(x-y) = -\frac{1}{4\pi|\bar{x}-\bar{y}|}\delta[|\bar{x}-\bar{y}| - (x^0 - y^0)]\theta(x^0 - y^0)$$

$$\bar{h}_{ab}(t,\bar{x}) = 4G \int d^3y \frac{T_{ab}(t_R,\bar{y})}{|\bar{x}-\bar{y}|} \qquad t_R \equiv t - |\bar{x}-\bar{y}|$$

Fourier transform

$$\tilde{\bar{h}}_{ab}(\omega,\bar{x}) = 4G \int d^3y \frac{e^{i\omega|\bar{x}-\bar{y}|}}{|\bar{x}-\bar{y}|} \tilde{T}_{ab}(t_R,\bar{y})$$

Production of GWs (2)

$$\tilde{\bar{h}}_{ab}(\omega,\bar{x}) = 4G \int d^3y \frac{e^{i\omega|\bar{x}-\bar{y}|}}{|\bar{x}-\bar{y}|} \tilde{T}_{ab}(t_R,\bar{y})$$

Isolated source, far away from the observer and slowly moving

$$\frac{e^{i\omega|\bar{x}-\bar{y}|}}{|\bar{x}-\bar{y}|} \approx \frac{e^{i\omega r}}{r}$$

Quadrupole formula (Lorenz gauge)

$$\bar{h}_{ij}(t,\bar{x}) = \frac{2G}{r} \left. \frac{d^2[I_{ij}]}{dt^2} \right|_{t_R=t-r}$$

$$I_{ij}(t) = \int d^3y y_i y_j T^{00}$$

Quadrupole moment tensor

No dipole (unlike in EM): conservation of the total momentum

- GWs are weak!
- GW production is coherent

Production of GWs (3)

Gravitational radiation emitted by a binary star:

$$\bar{h}_{ij}(t,\bar{x}) \propto \frac{8GM}{r} \Omega^2 R^2$$

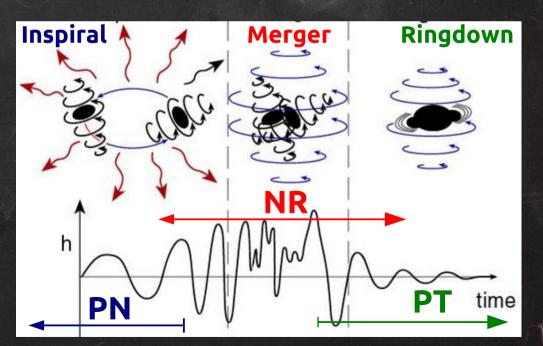
Massive object Large velocity Small distance from us

Energy loss in GWs $P=-\frac{32}{5}\frac{G^4M^3\mu^2}{c^5R^5}$ Quadrupole formula

In scalar-tensor theories:

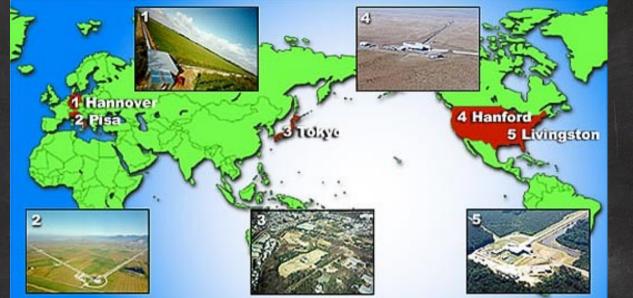
$$P = -\frac{2}{3} \frac{\mu^2 M^2}{R^4} \frac{\mathcal{S}^2}{\omega_{\rm BD}}$$

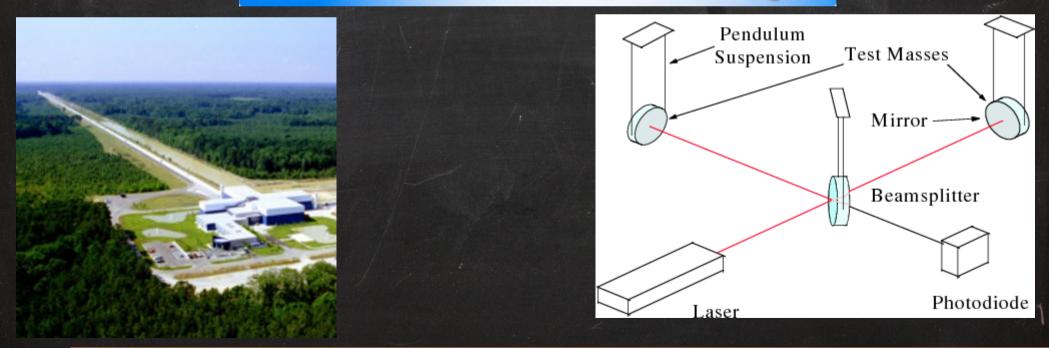
Dipole formula



Adapted from Thorne

GW detectors





Detection of GWs

$$f=\frac{\Omega}{2\pi}=\frac{cR_S^{1/2}}{4\pi\sqrt{2}R^{3/2}}$$
 frequency

$$\Omega = \left(\frac{GM}{4R^3}\right)^{T}$$
$$R_S = \frac{2GM}{c^2}$$

from Kepler's law

Schwarzschild radius

$$\bar{h}_{ij}(t,\bar{x}) \propto \frac{8GM}{r} \Omega^2 R^2 \Rightarrow h \sim \frac{R_S^2}{rR}$$

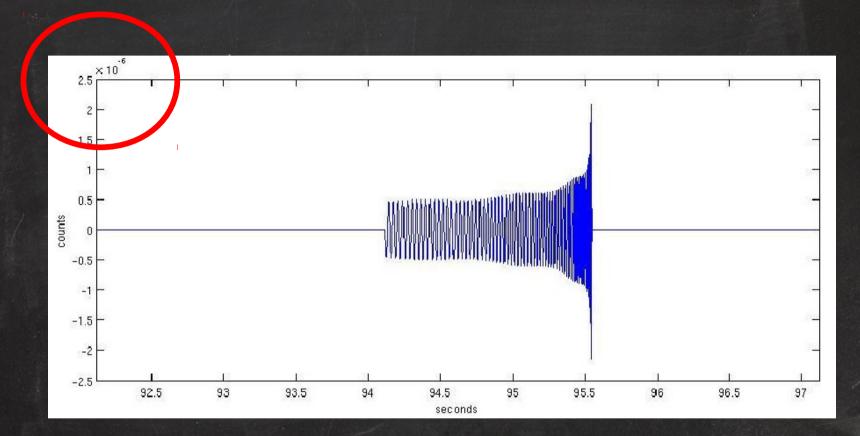
$$1\sim rac{\partial L}{L}$$
 sensitivit

/2

 $M \sim 10 M_{\odot}, \quad R \sim 10 R_S, \quad r \sim 100 \text{Mpc} \Rightarrow \begin{cases} f \sim 10^2 \text{Hz} \\ h \sim 10^{-21} \end{cases}$

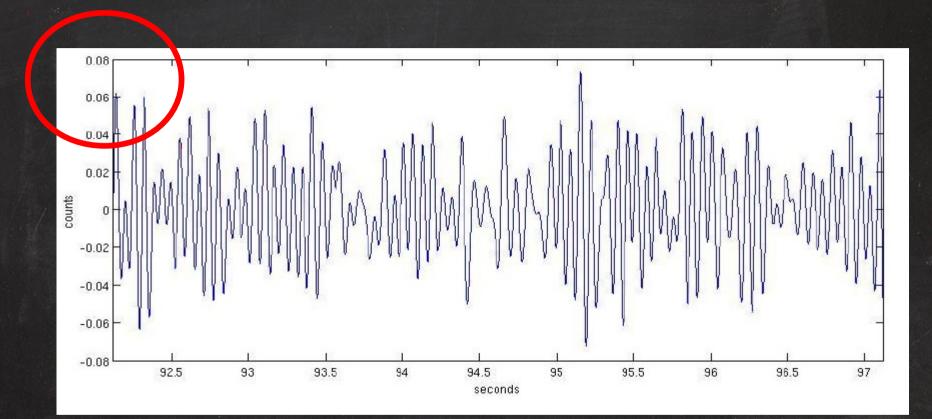
$$\delta L \sim 10^{-3} \frac{h}{10^{-21}} \frac{L}{\text{km}} \text{fm}$$

Detection of GWs



What you would like to see

Detection of GWs

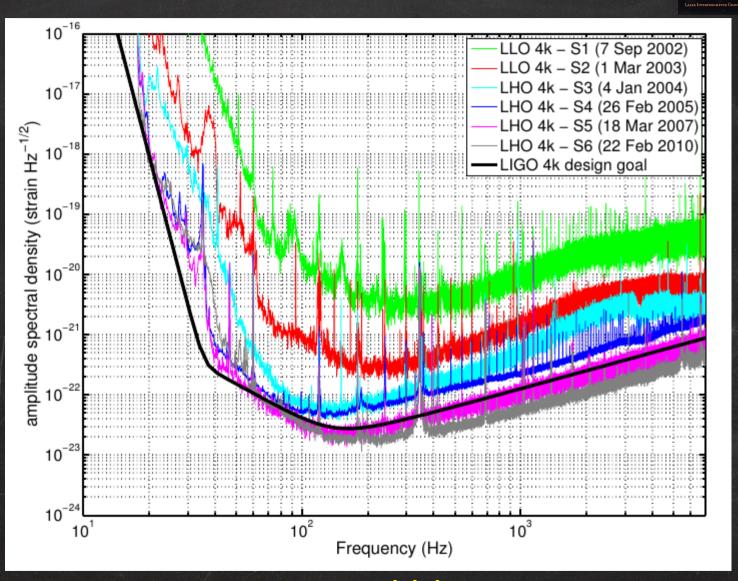


What you actually see

Extremely sensitive detector and

theoretical knowledge of the waveform

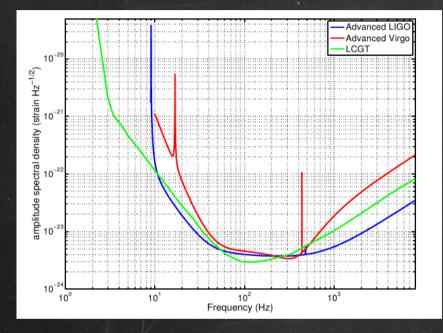
Sensitivity

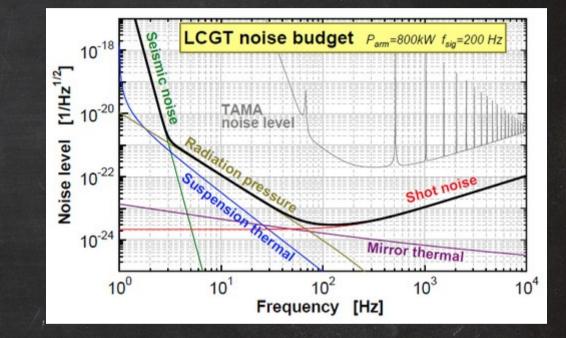


G

LIGO sensitivity Planes, trains, cars, earthquakes, ...

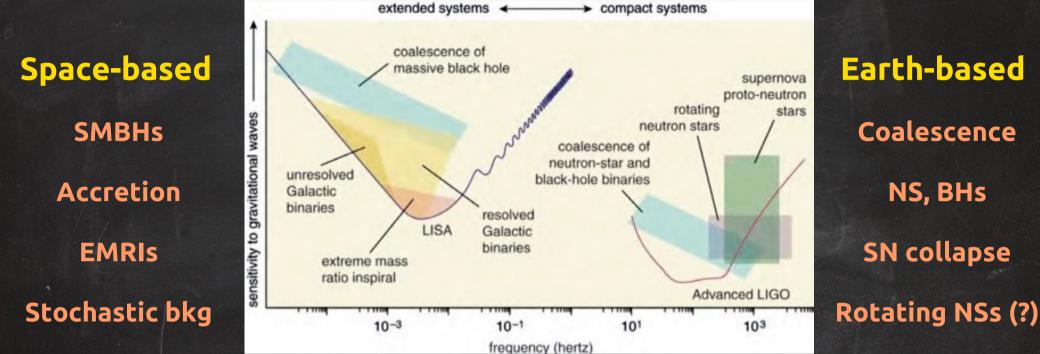
Second generation

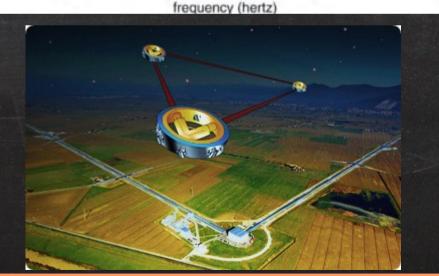




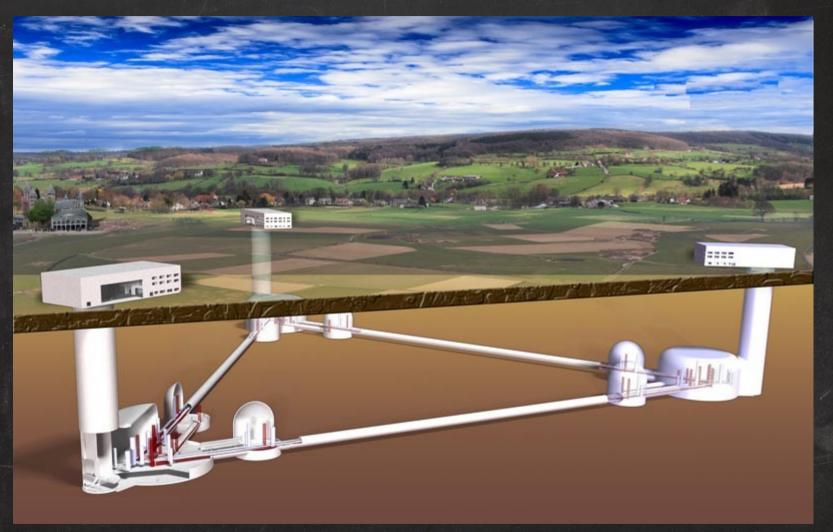
From 0.1 (at most!) to 40 NS-NS coalescences per yr

Future (?)





Future (?)



Einstein telescope ?

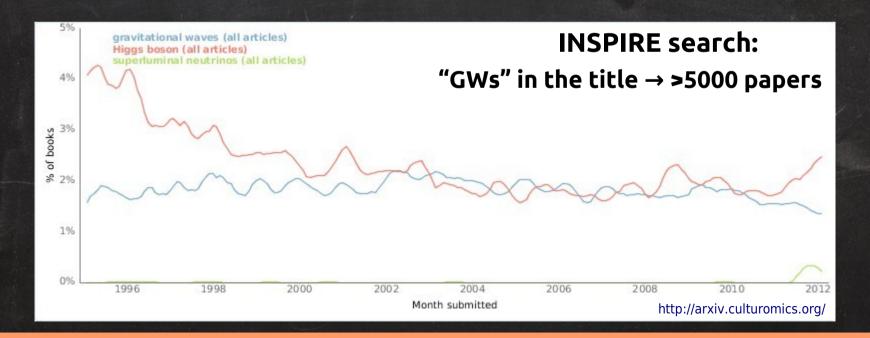
Conclusions

- GWs do exist!
- Their direct discovery will be one of the breakthroughs of this century
- Utmost importance for (astro)physics
- Tests of GR
- New era in astrophysics → GW astronomy
- New look at the whole universe



Some reference

- Textbooks: Hartle, MTW, D'Inverno, Carroll
- Reviews: Schutz & Ricci (gr-qc/1005.4735), Blair et al. (2012)
- LIGO reviews: Science 256 (1992) 325-333 and arXiv:0711.3041
- eLISA/NGO review: Amaro-Seoane et al. (arXiv:1201.3621)
- GW tests of GR: Will's review (gr-qc/0510072)



Want to detect GWs?



Join Einstein@home! http://einstein.phys.uwm.edu

einstein.

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Thanks!